

Phase separation and crystallisation of polydisperse hard spheres

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Hard spheres with a polydispersity above approximately 8% are shown to crystallise into two phase-separated solid phases. A polydispersity above 8% is too large to be tolerated by a single solid phase but phase separation produces two fractions with polydispersities sufficiently narrow to allow them to crystallise. It may not be possible to observe this in experiment due to the intervention of a glass transition.

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Experiments on synthetic colloids are inevitably on polydisperse colloids. The colloidal particles are not all the same, some are larger than average, some are smaller [1]. This has prompted considerable theoretical study of the phase behaviour of polydisperse systems, in particular the fluid–solid transition of polydisperse hard spheres [2–6]. The conclusion of these studies is that a solid of hard spheres can tolerate a polydispersity of a little less than 0.1, beyond this it is never stable with respect to the fluid phase. The 0.1 refers to the ratio of the standard deviation to the mean, denoted below by s . As Pusey [7] has noted, this makes sense if we note that the lattice constant of the solid is at most 10% larger than the diameter of the hard spheres. Hard spheres melt at a fraction $\simeq 0.74$ of close packing [8] which corresponds to a lattice constant $\simeq 1.11$ times the hard sphere diameter. Only if the polydispersity is sufficiently narrow that few or no spheres are larger than this lattice spacing can the mixture crystallise [7]. We have derived an analytical free energy for the solid phase of polydisperse hard spheres and our results agree with those of previous authors [2–6]. We find that the Helmholtz free energy of the solid phase of polydisperse hard spheres with $s > 0.082$ is never lower than that of the fluid phase at the same density.

However, all this takes no account of possible demixing. If the spheres phase separate into fractions then these fractions will have narrower polydispersities than the original mixture and so can crystallise. Of course, this phase separation costs some ideal mixing entropy but at high density spheres pack more efficiently in the solid phase than in the fluid resulting in a larger entropy in the solid. The possibility of demixing has been mentioned before [2,7] but no calculations have been performed. Here,

calculations show polydisperse hard spheres crystallising into two solid phases, one containing the larger spheres and one containing the smaller spheres. There is no fluid–fluid phase separation at the moderate polydispersities we consider, although there may well be at much larger polydispersities [9,10]. By moderate we mean polydispersities up to about 20%.

The crystallisation of polydisperse hard-sphere-like colloids into more than one solid phase has not been observed experimentally. This is despite the fact that it is almost certainly the equilibrium behaviour. The discrepancy between theory and experiment may be due to the intervention of the glass transition [11,12]. The cost in demixing entropy and the less efficient packing of polydisperse spheres into a solid vis-a-vis monodisperse spheres both act to raise the coexisting densities of the fluid–solid transition. It is quite possible that when the polydispersity is sufficiently large for the equilibrium behaviour to be crystallisation into two solid phases, the transition occurs at densities above the glass transition and is thus impossible to observe in experiment. Already for near-hard-sphere colloids with a polydispersity $s = 0.075$ the fluid–solid and glass transitions appear to be very close [1].

In a polydisperse mixture of spheres, spheres with a range of diameters are present [13,14]; there is a continuous distribution of sizes of spheres with a number density $\rho x(\sigma)d\sigma$ of spheres of size σ . ρ is the total number density of spheres. The width of the polydispersity is characterised by a width parameter w . The larger w is the broader the distribution of sizes present in the mixture. In the limit $w \rightarrow 0$ we recover a monodisperse system. The functional form of the polydispersity should not matter too much; we select a simple form, the hat

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function:

$$x(\sigma) = \begin{cases} 0 & \sigma < \bar{\sigma} - w/2 \\ w^{-1} & \bar{\sigma} - w/2 \leq \sigma \leq \bar{\sigma} + w/2 \\ 0 & \sigma > \bar{\sigma} + w/2 \end{cases}, \quad (1)$$

where $\bar{\sigma}$ is the mean diameter. The standard deviation divided by its mean, of the distribution $x(\sigma)$ is $s = w/(\sqrt{12}\sigma)$.

An approximation to the Helmholtz free energy of the solid phase may be obtained using a cell theory [15]. Our system is athermal so the free energy is equal to minus the entropy (ignoring the contribution from the momenta). Within a cell theory the entropy per particle is just the logarithm of the configuration space available to one particle trapped in a cell formed by its neighbours. These neighbouring particles are taken to be fixed at their lattice positions, which are the lattice constant a from the centre of the cell. For a solid of monodisperse spheres of diameter σ the hard-core repulsion prevents the particle from moving more than $(a - \sigma)$ from its lattice position in the direction of one of its neighbours. Then we can approximate the cell by a sphere of radius $a - \sigma$ [15] and so obtain a free energy per particle f_s in the solid phase

$$\frac{f_s}{kT} = -\ln \left[\frac{4\pi}{3} (a - \sigma)^3 \right], \quad (2)$$

where for a close-packed lattice, $a = (\eta_{cp}/\eta)^{1/3}$; k is Boltzmann's constant and T is the temperature. The volume fraction η of monodisperse spheres is related to the number density by $\eta = (\pi/6)\rho\sigma^3$, and its maximum value is $\eta_{cp} = \pi/(3\sqrt{2})$, which is achieved at close packing.

In order to generalise Eq. (2) to a solid of polydisperse spheres we have to make a further approximation. We neglect any correlations between the size of spheres at adjacent lattice points. In reality we would expect that on average smaller spheres will tend to cluster around large spheres to reduce the strain on the lattice. Every sphere is assumed to be in a cell formed from spheres of the average diameter $\bar{\sigma}$. Then the cell of a sphere of size σ is of radius $(a - (\sigma + \bar{\sigma})/2)$. This gives us the free energy of a sphere of size σ ; we simply integrate over all sizes of spheres, weighting with $x(\sigma)$ to obtain the free energy per particle

$$\begin{aligned} \frac{f_s}{kT} &= -\ln(4\pi/3) - 3 \int x(\sigma) \ln[a - (\sigma + \bar{\sigma})/2] d\sigma - \ln w \\ &= -\ln(4\pi/3) + 3 \left[1 + \frac{2}{w} (a - \bar{\sigma} - w/4) \ln(a - \bar{\sigma} - w/4) \right. \\ &\quad \left. - \frac{2}{w} (a - \bar{\sigma} + w/4) \ln(a - \bar{\sigma} + w/4) \right] - \ln w, \end{aligned}$$

where the last term is the ideal mixing entropy [13] for our polydisperse mixture [13,14]. The lattice constant a is still derived from $a = (\eta_{cp}/\eta)^{1/3}$ with η_{cp} for a

monodisperse solid. For polydisperse spheres with a distribution Eq. (1) the volume fraction is related to the total number density by $\eta = (\pi/6)\rho(1 + w^2/4)$. For the free energy of the fluid phase we use the polydisperse generalisation of Salacuse and Stell [13] of the free energy of Boublik, and Mansoori, Carnahan, Starling and Leland (BMCSL) [16,17].

Finding phase coexistence for polydisperse mixtures is very involved [2,3,13,14], but see Refs. [18,19]. Determining whether the fluid or the solid phase is more stable, has the lower free energy, at a given density is however straightforward, so we will content ourselves with this. In Fig. 1 we have plotted the free energies of the fluid and solid phases with a polydispersity $w/\bar{\sigma} = 0.3$ ($s = 0.087$). The free energy of the solid phase is always higher than that of the fluid phase and so is never the equilibrium phase. This holds if and only if the polydispersity is greater than $w = 0.28$ ($s > 0.082$). However, the free energy of a solid phase with half the width, here $w/\bar{\sigma} = 0.15$ does drop below that of the fluid phase. As the fluid has a higher free energy than solid phases with half the polydispersity it is unstable with respect to them. The phase separation will not of course be that complete or simple. Describing the two phases as one with all the spheres smaller than $\bar{\sigma}$ and the other with the spheres larger than $\bar{\sigma}$ is a gross oversimplification. However, we have proven that the fluid becomes unstable with respect to two solid phases, whilst being stable with respect to one. We have also plotted the free energy of the solid phase with $w/\bar{\sigma} = 0.1$. In the range of densities plotted in Fig. 1 it is always higher than for $w/\bar{\sigma} = 0.15$ so we expect that on crystallisation phase separation into two not three solid phases will occur. Also, the free energy of the fluid phase with $w/\bar{\sigma} = 0.3$ is lower than that of the fluid phase for either $w/\bar{\sigma} = 0.15$ or 0.1 . As an aside, we note that polydisperse sticky spheres also phase separate into more than one solid phase [20].

The free energy of phase separated solid phases only drops below that of the fluid phase above a volume fraction $\eta = 0.6$. Experiments on near-hard-sphere colloids with a polydispersity $s \approx 0.04$ [12] show a glass transition at $\eta = 0.58$. Although of course we have not been able to find the coexisting fluid and solid phases this at least suggests that, in experiment, the combined phase-separation-crystallisation phase transition may be prevented by a glass transition. Even if it is not, the dynamics of a combined phase separation and crystallisation will be slow [7,21].

Hard spheres with a polydispersity width $w > 0.28$ ($s > 0.082$) cannot crystallise into a single solid phase, its free energy is always higher than that of the fluid. (3) However, they can crystallise into two solid phases, we have verified this for polydispersities up to $w = 0.4$ ($s = 0.12$). We have not performed calculations beyond this because our solid phase free energy may not be reliable at these polydispersities, and because the density

at which the fluid and solid free energies are equal are so high ($\eta > 0.63$) at these polydispersities that the BMCSL free energy of the fluid phase may no longer be reliable.

In conclusion, experiments on near-hard-sphere colloids only find crystallisation when the polydispersity is no more than approximately $s = 0.075$ [7,1]. This is not the equilibrium behaviour of polydisperse hard spheres. At equilibrium, as the polydispersity is increased a second solid phase appears; we predict that this occurs at a polydispersity $s < 0.082$. Each of the two solid phases is about half as polydisperse as the original fluid. For polydispersities s of tens of per cent the picture is much less clear: there will now be spheres with very different diameters present which may favour the formation of lattice types other than a substitutionally-disordered close-packed lattice. The formation of very different lattices has been observed for binary mixtures of hard spheres [22,23]. Another possibility is that highly polydisperse hard spheres do not have an equilibrium solid phase.

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- [5] BARTLETT P., *J. Chem. Phys.*, **107** (1997) 188.
- [6] BARTLETT P., and WARREN P. B., preprint.
- [7] PUSEY P. N., *J. Phys. (France)*, **48** (1987) 709.
- [8] HOOVER W. G., and REE F. H., *J. Chem. Phys.*, **49** (1968) 3609.
- [9] CUESTA J. A., preprint.
- [10] WARREN P. B., preprint.
- [11] VAN MEGEN W. and PUSEY P. N., *Phys. Rev. A*, **43** (1991) 5429.
- [12] UNDERWOOD S. M. and VAN MEGEN W., *Phys. Rev. E*, **49** (1994) 4206.
- [13] SALACUSE J. J., and STELL G., *J. Chem. Phys.*, **77** (1982) 3714.
- [14] GUALTIERI J. A., KINCAID J. M., and MORRISON G., *J. Chem. Phys.*, **77** (1982) 521.
- [15] BUEHLER R. J., WENTORF R. H., HIRSCHFELDER J. O., and CURTISS C., *J. Chem. Phys.*, **19** (1951) 61.
- [16] BOUBLIK T., *J. Chem. Phys.* **53** (1970) 471.
- [17] MANSOORI G. A., CARNAHAN N. F., STARLING K. E. and LELAND T. W., *J. Chem. Phys.* **54** (1971) 1523.
- [18] SOLLICH P., and CATES M. E., *Phys. Rev. Lett.*, **80** (1998) 1365.
- [19] WARREN P. B., *Phys. Rev. Lett.*, **80** (1998) 1369.
- [20] SEAR R. P., cond-mat/9805201 (<http://xxx.lanl.gov/>).
- [21] HENDERSON S. I. and VAN MEGEN W., *Phys. Rev. Lett.*, **80** (1998) 877.
- [22] BARTLETT P., OTTEWILL R. H. and PUSEY P. N., *Phys. Rev. Lett.*, **68** (1992) 3801.
- [23] ELDRIDGE M. D., MADDEN P. A. and FRENKEL D., *Nature*, **365** (1993) 35.

Fig. 1. The free energy of polydisperse hard spheres in the fluid and solid phase for several polydispersities. The solid curve is for the fluid phase with a polydispersity width $w = 0.3$. The dotted, dashed and dot-dashed curves are for the solid phase with polydispersity widths $w = 0.3, 0.15$ and 0.1 , respectively. The dotted curve has been cut off just before a rapid divergence.

